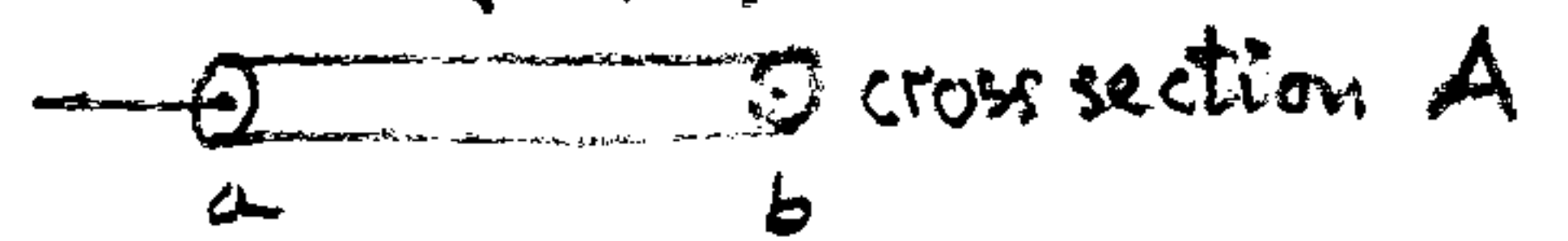


Finite Volume Discretization of  $\frac{\partial u}{\partial t} + \frac{\partial F}{\partial x} = 0$  in  $\Omega = (a, b)$



mesh:  $x_1, \dots, x_{i-1}, x_i, \dots, x_M$   
 nodes  $x_i, i=1:M$  (and  $x_0=a, x_{M+1}=b$ )  
 faces  $x_{i-\frac{1}{2}}, i=1:M+1$   
 control volumes  $V_i = [x_{i-\frac{1}{2}}, x_{i+\frac{1}{2}}] \times A$   
 $\Delta V_i = A \cdot \Delta x$

Integrate over  $V_i$  and over  $[t_n, t_{n+1}]$ :

$$\int_{t_n}^{t_{n+1}} \int_{V_i} \frac{\partial}{\partial t} u(x,t) A dx dt + \int_{t_n}^{t_{n+1}} \left[ AF(x_{i+\frac{1}{2}}, t) - AF(x_{i-\frac{1}{2}}, t) \right] A dx dt = 0$$

$$\Rightarrow \int_{V_i} u A dx \Big|_{t=t_n}^{t=t_{n+1}} + \int_{t_n}^{t_{n+1}} [AF(x_{i+\frac{1}{2}}, t) - AF(x_{i-\frac{1}{2}}, t)] dt = 0$$

Set  $U_i^n := \frac{1}{\Delta V_i} \int_{V_i} u(x, t_n) dV = \text{mean value of } u \text{ over } V_i \text{ at time } t_n$

$AF_{i-\frac{1}{2}}^{n+\theta} := \frac{1}{\Delta t} \int_{t_n}^{t_{n+\theta}} AF(x_{i-\frac{1}{2}}, t) dt = \text{mean flow rate across face } i-\frac{1}{2} \text{ during time step}$

$$\Rightarrow \Delta V_i [U_i^{n+1} - U_i^n] + \Delta t [AF_{i+\frac{1}{2}}^{n+\theta} - AF_{i-\frac{1}{2}}^{n+\theta}] = 0 \quad \begin{matrix} i=1:M \\ n=0,1,\dots,N_{\text{steps}} \end{matrix}$$

expressing exact discrete conservation over each Control Volume and each time-step.

Explicit scheme: choose  $\theta=0$ : assume mean flux during  $\Delta t$  remains  $\approx$  flux at time  $t_n$ .

Knowing  $U_i^n$  and flux law  $F=F(u)$ , we can compute  $F_{i-\frac{1}{2}}^n, i=2:M$

and BCs provide  $F_{1/2}^n, F_{M+1/2}^n$ , so we can update

$$U_i^{n+1} = U_i^n + \frac{\Delta t}{\Delta V_i} [AF_{i-\frac{1}{2}}^n - AF_{i+\frac{1}{2}}^n], \quad i=1:M, n=0,1,\dots,N$$

Implicit schemes: choose  $0 < \theta \leq 1$  and  $F^{n+\theta} = (1-\theta)F^n + \theta F^{n+1}$

Then we have a coupled system to solve for  $\{U_i^{n+1}\}$  at each timestep (using some system solver)

$\theta = \frac{1}{2}$  is Crank-Nicolson scheme

$\theta = 1$  is Fully Implicit (Backward Euler) scheme

Note: Cross-sectional area  $A$  cancels out in 1D, but not in 2D, 3D

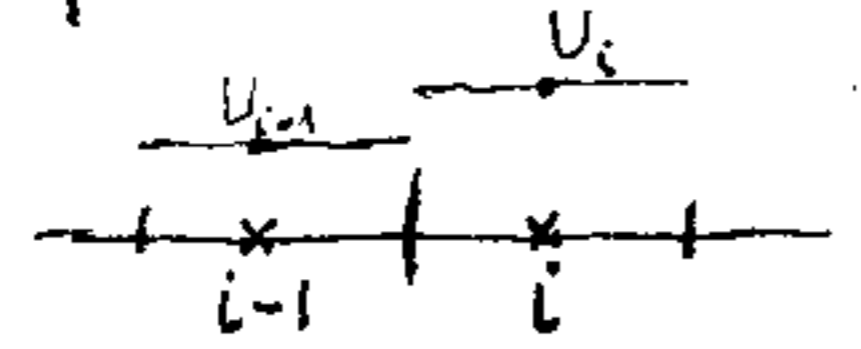
# Flux for diffusion / heat conduction

The conservation law  $u_t + F_x = S$  conserves  $u$  for any  $F$

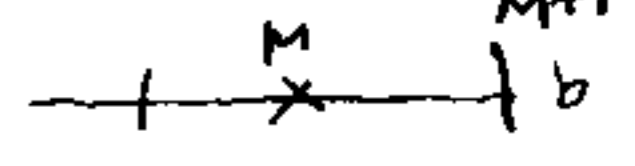
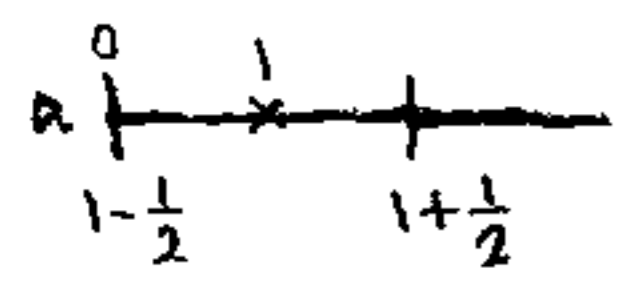
$F$  is specified by a constitutive law expressing particular physical phenomena.

Diffusion: Fick's law:  $F = -D u_x$ ,  $D = \text{diffusivity [cm}^2/\text{sec]}$   
 $u = \text{concentration [gr/cm}^3]$

( $\Rightarrow$  diffusion equation:  $u_t = (D u_x)_x = D u_{xx}$  if  $D \equiv \text{constant}$   
 a most important PDE, more later...)



Flux Discretization:  $F_{i-\frac{1}{2}} \approx -D \frac{u_i - u_{i-1}}{\Delta x}$ ,  $i = 2 : M$  (internal faces)  
 $F_{1-\frac{1}{2}} \approx -D \frac{u_1 - u_0}{\Delta x/2}$ ,  $F_{M+\frac{1}{2}} \approx -D \frac{u_{M+1} - u_M}{\Delta x/2}$   
 boundary fluxes  $F_{\frac{1}{2}}$ ,  $F_{M+\frac{1}{2}}$  from Boundary Conditions



e.g. impermeable boundary means Flux = 0 (so nothing goes through the bry)  
 so  $F_{\frac{1}{2}} = 0$  and  $F_{M+\frac{1}{2}} = 0$

# Boundary Conditions in FV scheme for diffusion

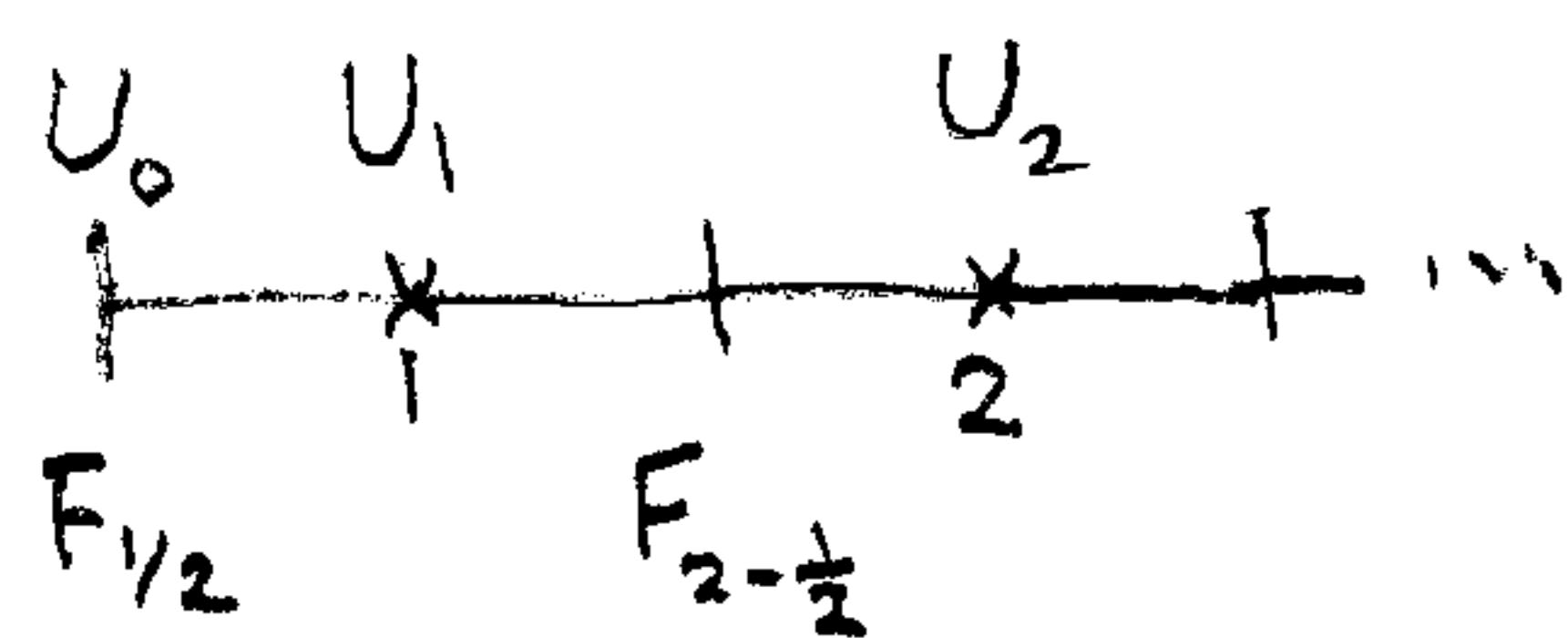
A PDE Problem consists of: PDE e.g.  $u_t = D u_{xx}$ ,  $x \in \Omega$ ,  $t_0 < t < t_{end}$   
 for  $u(x, t)$  IC  $u(x, 0) = u_{init}(x)$ ,  $x \in \Omega$

Scheme needs fluxes

BCs some BC at each pt of  $\partial\Omega$  for all time

Standard BCs at  $x=a$  or  $x=b$  coded in FLUX routine

I. Dirichlet (or 1<sup>st</sup> kind): impose value at  $x=a$ :  $u(a, t) = u_a(t)$  given



so  $U_0^n \approx u(a, t_n) = u_a(t_n)$  known by value

$$\Rightarrow F_{1/2}^n = -D \frac{U_1^n - U_0^n}{\Delta x/2}, \quad F_{m+1/2}^n = -D \frac{U_{m+1}^n - U_m^n}{\Delta x/2}$$

II. Neumann (or 2<sup>nd</sup> kind): impose flux at  $x=a$ :  $\mathcal{F}(a, t) = F_a(t)$  given

so  $F_{1/2}^n = F_a(t_n)$  known by flux

then by value  $U_0^n$  can be found from  $F_{1/2} = -D \frac{U_1 - U_0}{\Delta x/2} = F_a(t_n)$

$$\Rightarrow U_0^n = U_1^n + \frac{\Delta x}{2D} F_{1/2}^n, \quad U_{m+1}^n = U_m^n - \frac{\Delta x}{2D} F_{m+1/2}^n$$

solve for  $U_0$

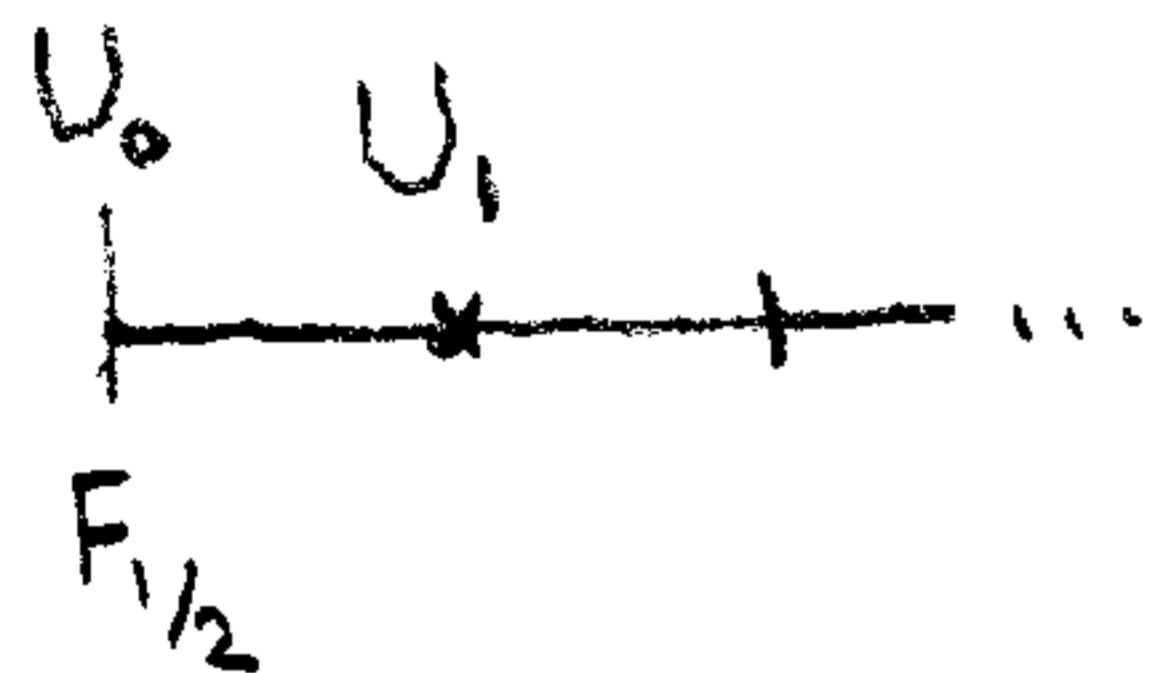
In particular, impermeable/insulated by:  $\mathcal{F}(a, t) = F_a(t) \equiv 0$

$$\Rightarrow U_0^n = U_1^n, \quad U_{m+1}^n = U_m^n$$

Note: FV scheme needs only fluxes. Boundary values  $U_0, U_{m+1}$  (consistent with BCs) are for outputting (to plot profiles).

BCs ...

III. Convective BC (or Robin or 3<sup>rd</sup> kind):  
most realistic BC physically



$$f(a, t) = h [u_{\text{amb}}(t) - u(a, t)]$$

$h$  = film coeff,  
heat transfer coeff,

$u_{\text{amb}}(t)$  = ambient  $u$ , given

$$F_{1/2}^n = h [u_{\text{amb}}(t_n) - U_0^n] \stackrel{\text{want}}{=} -D \frac{U_1^n - U_0^n}{\Delta x/2}$$

solve for  $U_0^n$ :

$$U_0^n = \frac{U_1 + \frac{h \Delta x}{2D} u_{\text{amb}}(t_n)}{1 + \frac{h \Delta x}{2D}}$$

for Output

$$\Rightarrow F_{1/2}^n = -D \frac{U_1^n - u_{\text{amb}}(t_n)}{\frac{D}{h} + \frac{\Delta x}{2}}$$

so  $u_{\text{amb}}$  plays role of  $U_0$   
but with  $\frac{D}{h} + \frac{\Delta x}{2}$  instead of  $\frac{\Delta x}{2}$

Note:  $h \rightarrow \infty \Rightarrow u(a, t) = u_{\text{amb}}(t)$  (for  $f$  to be finite).  
Dirichlet BC!

$h = 0 \Rightarrow F = 0$ , insulated by  
homogeneous Neumann BC