Pattern Formation During Directional Epitaxy

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Epitaxial Growth and Continuous Processing

- Epitaxy describes ordered crystalline growth on a mono-crystalline substrate following the crystal structure of the substrate.
- Continuous processing has been used to create large quantities of thin film coated tapes/wire using a reel-to-reel system [Cui et al.,IEEE Trans. Appl. Supercond.,1999].
- "Directional Epitaxy" [Schulze, J. Crystal Growth, 2006] describes epitaxy on a continuously supplied substrate.
- For small deposition zones or near the ends of a larger system, finite size effects and boundary conditions are important.



1-D Continuum Model



Geometry: *n* steps, n+1 terraces (side view)

1-D Burton-Cabrera-Frank (BCF) continuum model [Burton et al., Phil. Trans. Roy. Soc., 1951]:

$$\begin{split} \partial_t \rho_j - V \partial_x \rho_j &= D \partial_x^2 \rho_j + F, \quad \xi_j < x < \xi_{j+1}, \ j = 0, \dots, n \\ \pm D \partial_x \rho \big|_{\pm} &\pm (V + \partial_t \xi_j) \rho \big|_{\pm} = k_{\pm} (\rho - \rho_e) \big|_{\pm}, \quad x = \xi_j(t), \ j = 1, \dots, n \\ \rho_a(V + \partial_t \xi_j) &= D [\partial_x \rho]_-^+ + (V + \partial_t \xi_j) [\rho]_-^+, \quad x = \xi_j(t), \ j = 1, \dots, n \\ \rho_n &= \rho_0 = \rho_e, \quad x = \pm L, \text{ respectively} \end{split}$$

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1-D BCF Model

$$\begin{split} \partial_t \rho_j - V \partial_x \rho_j &= D \partial_x^2 \rho_j + F, \quad \xi_j < x < \xi_{j+1}, \ j = 0, \dots, n \\ \pm D \partial_x \rho \big|_{\pm} &\pm (V + \partial_t \xi_j) \rho \big|_{\pm} = k_{\pm} (\rho - \rho_e) \big|_{\pm}, \quad x = \xi_j(t), \ j = 1, \dots, n \\ \rho_a(V + \partial_t \xi_j) &= D [\partial_x \rho]_-^+ + (V + \partial_t \xi_j) [\rho]_-^+, \quad x = \xi_j(t), \ j = 1, \dots, n \\ \rho_n &= \rho_0 = \rho_e, \quad x = \pm L, \text{ respectively} \end{split}$$

Saum, Schulze (Dept of Mathematics, UTK) Directional Epitaxy Pattern Formation

LIT.

 Quasistatic approximation assumes the adatom density equilibrates fast compared to the motion of the steps.

$$\begin{split} &\partial_t \rho_j - V \partial_x \rho_j = D \partial_x^2 \rho_j + F, \quad \xi_j < x < \xi_{j+1}, \ j = 0, \dots, n \\ &\pm D \partial_x \rho \big|_{\pm} \pm (V + \partial_t \xi_j) \rho \big|_{\pm} = k_{\pm} (\rho - \rho_e) \, \big|_{\pm}, \quad x = \xi_j(t), \ j = 1, \dots, n \\ &\rho_a(V + \partial_t \xi_j) = D [\partial_x \rho]_-^+ + (V + \partial_t \xi_j) [\rho]_-^+, \quad x = \xi_j(t), \ j = 1, \dots, n \\ &\rho_n = \rho_0 = \rho_e, \quad x = \pm L, \text{ respectively} \end{split}$$

L1r

- Quasistatic approximation assumes the adatom density equilibrates fast compared to the motion of the steps.
- $\rho = \rho \rho_e$, scale by the diffusive time scale a^2/D

$$\begin{split} \mathbf{0} &= \mathbf{D}\partial_x^2 \rho_j + \mathbf{F}, \quad \xi_j < x < \xi_{j+1}, \ j = 0, \dots, n \\ &\pm \mathbf{D}\partial_x \rho \big|_{\pm} = \mathbf{k}_{\pm} \left(\rho - \rho_e \right) \big|_{\pm}, \quad x = \xi_j(t), \ j = 1, \dots, n \\ &\rho_a(\mathbf{V} + \partial_t \xi_j) = \mathbf{D}[\partial_x \rho]_-^+, \quad x = \xi_j(t), \ j = 1, \dots, n \\ &\rho_n = \rho_0 = \rho_e, \quad x = \pm L, \text{ respectively} \end{split}$$



L1r

- Quasistatic approximation assumes the adatom density equilibrates fast compared to the motion of the steps.
- $\rho = \rho \rho_e$, scale by the diffusive time scale a^2/D
- 1-D BCF Quasistatic Approximation

$$\begin{split} 0 &= \partial_x^2 \rho_j + F, \quad \xi_j < x < \xi_{j+1}, \ j = 0, \dots, n \\ \partial_x \rho_j &= \rho_j, \quad x = \xi_j, \ j = 1, \dots, n \\ -\partial_x \rho_{j-1} &= \rho_{j-1}, \quad x = \xi_j, \ j = 1, \dots, n \\ V &+ \partial_t \xi_j = \partial_x \left(\rho_j - \rho_{j-1} \right), \quad x = \xi_j, \ j = 1, \dots, n \\ \rho_n &= \rho_0 = 0, \quad x = \pm L, \ \text{respectively.} \end{split}$$

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1-D Dynamical System

The adatom density takes the form:

$$\rho_j(\mathbf{x};t) = A_j(\vec{\xi};t) + B_j(\vec{\xi};t)\mathbf{x} - \frac{F\mathbf{x}^2}{2}$$

We then obtain the ODE system:

$$\partial_t \xi_j = B_j(\vec{\xi},t) - B_{j-1}(\vec{\xi},t) - V, \quad j = 1,\ldots,n.$$

where for fixed set of $\vec{\xi}$, one can easily solve for $\{A_j, B_j\}_{j=0}^n$ the following block diagonal linear system:

$$\mathbf{M}(\vec{\xi}) \begin{pmatrix} A_0 \\ B_0 \\ A_1 \\ B_1 \\ \vdots \\ A_n \\ B_n \end{pmatrix} = \vec{b}(\vec{\xi})$$

Pairwise Step Patterns (PSP)

- A Pairwise Step Pattern (PSP) is an alternating interior terrace width pattern denoted by δ_0 and δ_1 .
- Step locations ξ_j , j = 1, ..., n can be expressed as:

$$\xi_j = \begin{cases} -L + \delta_L + \frac{j}{2}\delta_0 + \left(\frac{j}{2} - 1\right)\delta_1, & j \text{ even} \\ -L + \delta_L + \frac{j-1}{2}\delta_0 + \frac{j-1}{2}\delta_1, & j \text{ odd} \end{cases} \qquad j = 1, \dots, n.$$

- BCF approximation indicates PSP's exist and have fundamentally different character when the number of steps n is odd or even.
- When seeking steady states, for a range of V(n), there exists fixed points.

Steady State Adatom Density Profiles $\rho(x)$ (side view)

- Note the existence of PSP for both *n* odd and *n* even.
- Note the asymmetry for *n* odd. The two fixed points have the end profiles interchanged.



$\Delta - V$ Relation

- For *n* even, a wide range of velocities produces a wide range of δ₀'s and the Jacobian evaluated at the valid fixed points has purely imaginary eigenvalues.
- For *n* odd, a narrow range of velocities produces two values of δ₀ for each *V* and the Jacobian evaluated at the two fixed points has eigenvalues, one with all negative real parts and one with all positive real parts.



Kinetic Monte Carlo (KMC)

- Atomistic growth model in 2+1 dimensions with nearest neighbor interactions.
- State of surface described by integer height h(x, y) on a square grid lattice of dimensions M × N.
- Equilibrium boundary conditions at grid boundaries in the *x* direction; periodic boundary conditions in the *y* direction.
- We utilize an in-plane lateral nearest neighbor bond counting technique to determine rates (probabilities) of an adatom moving [Clarke and Vvedensky, J. Appl. Phys., 1988].
- Surface atoms hop to neighboring sites with rates given by *R*:

$$\Delta E = E_s + nE_n$$
, $R = K(T) \exp(-\Delta E/(k_B T))$

These hopping rates vary based on the current surface morphology.

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KMC Step Edge Contours (plan view)

- Evidence of PSP in 2+1 KMC clearly visible.
- Requires relatively short terrace widths to reduce island nucleation ٠ influence.



V = 10.50, n = 14

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2+1 KMC

KMC $\bar{\delta_0}$, BCF δ_0

- KMC appears to match slope correctly for *n* even.
- Possible reasons for not matching include nucleation, integer resolution of KMC, other stochastic and 2-D effects.



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Adatom Kinetic Monte Carlo (AAKMC)

- 1-D version of KMC, no nucleation, developed in order to be more like our BCF approximation.
- Allows us to then go to larger terrace widths and thus minimize integer resolution issues and stochastic noise.
- Rather than describing the state of the surface, one tracks the positions of step edges and adatom locations.
- Adatom movement is treated as a random walk on flat surface covering whole domain.
- Adatoms do not interact with each other.
- When adatom lands in front of step edge, that step edge moves forward.



AAKMC $\bar{\delta_0}$, BCF δ_0

• Note very close matching throught velocity range.



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Summary

- Nonperiodic BC are relevant for small systems or near the boundaries. Periodic BC are applicable to large systems or far from the boundaries.
- "Steady States" exist for continuously processed systems, i.e., the height profile is stationary and mean step positions don't change.
- These states take the form of a Pariwise Step Pattern (PSP) where the interior terrace widths alternate between two widths, δ₀ and δ₁.
- AAKMC confirms this with very good agreement with 1-D continuum approximation.

